

Evidence of structural phase transition in superconducting $\text{SmFeAsO}_{1-x}\text{F}_x$ from ^{19}F NMR

M. Majumder[†], K. Ghoshray^{††}, C. Mazumdar[†] A. Poddar[†] A. Ghoshray[†] D. Berardan[‡], N. Dragoe[‡]

[†] ECMP Division, Saha Institute of Nuclear Physics, 1/AF Bidhannagar, Kolkata-700064, India

[‡] Institut de Chimie Moléculaire et des Matériaux d'Orsay, Univ. Paris-Sud 11, UMR 8182, 91405 Orsay, France

Abstract.

We report the resistivity, magnetization and ^{19}F NMR results in polycrystalline sample of $\text{SmFeAsO}_{0.86}\text{F}_{0.14}$. The resistivity and the magnetization data show a sharp drop at 48 K indicating the superconducting transition. The nuclear spin-lattice rate ($1/T_1$) and spin-spin relaxation rate ($1/T_2$) clearly show the existence of a structural phase transition near 163 K in the sample which also undergo superconducting transition. This finding creates interest to explore whether this is unique for only Sm based systems or it is also present in other rare-earth based 1111 superconductors.

E-mail: kajal.ghoshray@saha.ac.in

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[†] To whom correspondence should be addressed (kajal.ghoshray@saha.ac.in)

1. Introduction

The existence of competition between structural, magnetic and superconducting (SC) transition [1] in F doped oxypnictides makes them unconventional. Such competing order parameters would lead to a complex phase diagram including several coexisting phases in the nano scale range. The general belief is that, the parent compound $ReFeAsO$ (Re = rare-earth) shows a structural phase transition (SPT) from tetragonal (T) to orthorhombic (O) symmetry around $T_s=160$ K and then nearly 20 K below T_s , there is a long range magnetic order (MO) with spin density wave (SDW) type transition. As F doping is increased, transition temperatures for SPT and MO gradually decrease and at a particular concentration of F, the system undergoes SC transition [2, 3].

Based on synchrotron powder x-ray diffraction and muon spin rotation studies, a new phase diagram for $SmFeAsO_{1-x}F_x$ system has been proposed indicating that the SPT (T - O) in the range 175-155 K survives up to optimal F doped sample [4]. The difficulty to determine the occurrence of the SPT is due to the decrease of O distortion with the increase of F content. SPT can be magnetically driven, relieving magnetic frustration resulting from AFM neighbor and next neighbor interactions between local Fe moments [5, 6]. It can also be related to the nematic ordering [7] defined as the spontaneously broken C4 rotational symmetry of the square lattice due to either AFM spin-fluctuations or orbital ordering as was found in cuprates [8] and in the 122 family [9]. Theoretical results also show that in 1111 family, the SPT is driven by spin-fluctuations or orbital ordering [10]. Martinelli *et al* [4] claimed that the survival of SPT even for the optimal doped superconducting sample is due to the orbital ordering mediated symmetry breaking force and not driven by magnetic fluctuations.

We intend to study $SmFeAsO_{0.86}F_{0.14}$ using nuclear magnetic resonance (NMR) which is a very useful microscopic tool to probe the local magnetic and structural properties. NMR is expected to detect any role of the magnetic fluctuation on the SPT. It may be mentioned that neither the signature of the SC transition nor SPT was detected from ^{19}F nuclear spin-lattice relaxation rate ($1/T_1$) in $SmFeAsO_{0.85}F_{0.15}$ [11].

2. Experimental

The polycrystalline sample with nominal composition $SmFeAsO_{0.86}F_{0.14}$ was prepared through a solid state reaction route using Sm, As, FeF_2 , Fe and Fe_2O_3 as starting materials. All handlings were made in an Ar-filled glove box with less than 1 ppm O_2 and H_2O . SmAs alloy was first obtained by heating Sm and As under pure Ar in a closed silica tube at 900 °C during 12 h. The single phase nature of the alloy was confirmed by X-ray diffraction. SmAs was then carefully ground and sieved to less than 100 μm , and mixed in stoichiometric ratio with Fe, Fe_2O_3 and FeF_2 , and the resulting powder was pressed into $2\times 3\times 12$ mm³ bars under 200 MPa. These bars were heated two times with intermediate grinding and pressing, at 1150 °C during 48 h under argon in alumina crucibles sealed in closed silica tubes, with Ta pieces as getter. XRD patterns (figure 1)

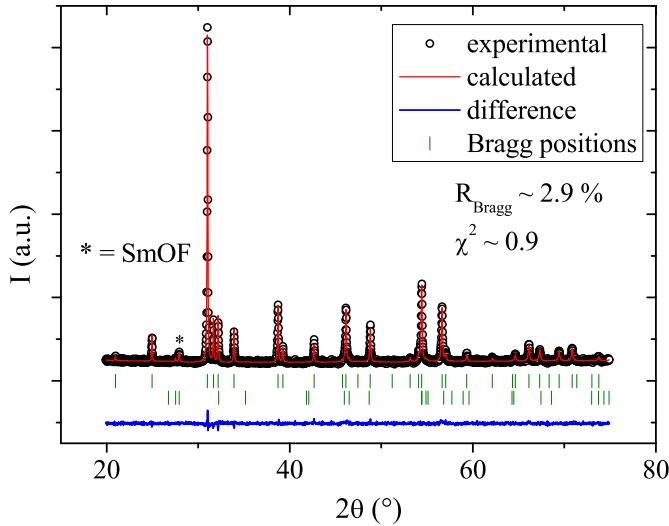


Figure 1. (Color online). XRD patterns of $SmFeAsO_{0.86}F_{0.14}$. Presence of faint amount of SmOF is marked by *. $\chi^2 = \sum_{i=1}^n w_i \{y_i - y_{ci}(\alpha)\}^2$ that is minimized in the Rietveld method [12], where y_i = profile intensity, y_{ci} = calculated counts with α as the parameter vector and $w_i = \frac{1}{\sigma_i^2}$, σ_i^2 being the variance of the y_i . The quality of the agreement between the observed and calculated profiles is measured by profile factor R_{Bragg} and is defined as $R_{Bragg} = 100 \frac{\sum_{i=1,n} |y_i - y_{ci}|}{\sum_{i=1,n} y_i}$.

were recorded using a Philips Xpert Pro diffractometer with XCelerator detector using $CuK\alpha 1$ radiation in a 2θ range $20 - 80^\circ$. The XRD patterns were analyzed using the Rietveld method with the help of the FullProf software [12].

It is very hard to determine the actual F content of 1111 compounds precisely, as F and O can not be distinguished using XRD and as the quantification of these light elements by EDX is challenging. However, we can say that the T_c is consistent [13] with a doping level of 0.14 so do the lattice parameters, obtained from Rietveld refinement, namely $a=3.9289(5)\text{\AA}$ and $c=8.467(1)\text{\AA}$ which are also consistent with this doping level as compared to the literature values [2].

The resistivity was measured using four probe method by applying magnetic fields in the range 0-8 T using a cryogen free magnet from M/S Cryogenics Limited. The temperature variation study was performed using a close cycle refrigerator also from M/S Cryogenics Limited. The DC magnetization measurements were done in the range 0.1-7 T in PPMS of M/S Quantum Design, Inc., USA.

The ^{19}F NMR measurements were performed on the powder sample in a fluorine free probe and carried out using a conventional phase-coherent spectrometer (Thamway

PROT 4103MR) with $H_0 = 7.04$ T superconducting magnet (Bruker). The spectrum was recorded by changing the frequency step by step and recording the spin echo intensity by applying a $\pi/2-\tau-\pi/2$ solid echo pulse sequence. The spin-lattice relaxation time (T_1) was measured using the saturation recovery method, applying a single $\pi/2$ pulse. The spin-spin relaxation time (T_2) was measured applying $\pi/2-\tau-\pi$ pulse sequence.

3. Results and discussion

3.1. Resistivity and magnetic susceptibility

Figure 2 shows the T dependence of resistivity (ρ) and the dc magnetic susceptibility (χ). $\rho(T)$ at zero field shows a clear signature of SC transition (T_c) at 48 K. The width of the transition (ΔT_c) = ± 0.8 K indicates the high homogeneity of the sample. Inset of figure 2a shows $\rho(T)$ at different magnetic field. As expected, ΔT_c is not as sharp as in zero field case. ρ , above T_c , does not vary with T^2 , usually seen in Fermi liquid case, instead it follows a linear behavior between T_c and 160 K. Hess *et al* [14] ascribed the linear behavior of $\rho(T)$ above T_c observed in $SmFeAsO_{1-x}F_x$ with x up to 0.1, as a strange metallic behavior.

Figure 2(b) shows variation of χ with T at 7 T with inset showing T dependence of M at different magnetic field. Well defined SC transition was observed at 48 K, even at 7 T. At $H \geq 1$ T, M is positive below T_c . This is because with the increase of field, the paramagnetic moment of Sm 4f (and/or Fe 3d) overcomes the diamagnetic response of the SC electrons.

$\chi - T$ curve, in the range 160 - 300 K, follows $\chi = \chi_0 + C/(T + \theta)$, with $\chi_0 = 0.0019$ emu/mole and $\theta = -11$ K. Where χ_0 contains the contributions from Pauli-paramagnetism, orbital paramagnetism and Landau diamagnetism of conduction electrons. From the Curie-Wiess (CW) constant (C), we found $P_{eff} = 0.79 \mu_B$, which is close to the free ion value ($0.84\mu_B$) of Sm^{3+} 4f local moment. The deviation of $\chi(T)$ from CW law may arise due to the development of short range magnetic correlation within Fe-3d spins. Contribution from Sm-4f is ruled out, as Sm 4f AFM ordering temperature being $T_N = 3.8$ K (manifested as a peak in the $\chi - T$ curve, shown in figure 2(b)), its contribution to χ is expected to follow the CW law near 160 K. This deviation may also result from the coupling between the Fe 3d spins/orbitals and the SPT as observed by Martinelli *et al* [4]. It is to be noted that the only secondary phase we have observed in the sample is from faint amount of SmOF, which is paramagnetic in the whole temperature range (see figure 1).

3.2. Nuclear magnetic resonance

Figure 3(a) shows ^{19}F NMR spectra along with the calculated one at different temperatures. Below 10 K, the spectra could not be detected because of the excessive line broadening/shortening of the nuclear spin-spin and spin-lattice relaxation times, due to the development of short range correlations among the Sm-4f spins as revealed

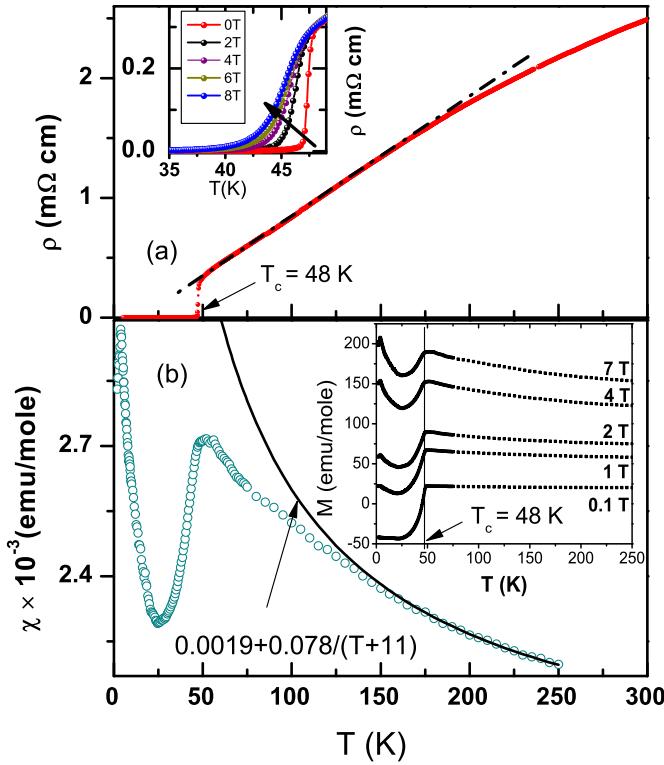


Figure 2. (Color online). Temperature dependence of resistivity (ρ) and susceptibility (χ) in $SmFeAsO_{0.86}F_{0.14}$. (a): $\rho(T)$ plot at zero field showing a sharp drop at $T_c = 48$ K, — · —: linear fit. Inset shows ρ at different magnetic field. (b): $\chi(T)$ at 7 T and the solid line is the Curie-Weiss fit. The inset shows magnetization (M) at different field.

in the magnetic susceptibility behavior. Almost symmetric (lorentzian) spectrum indicates negligible anisotropy in the local magnetic field at the ^{19}F site as was also reported in $(La/Sm)FeAsO_{1-x}F_x$ [15, 11]. The shift (K) of the resonance line (from the diamagnetic reference position) was measured either from the peak position of the spectra or obtained from the theoretically fitted curve and is negative in sign (figure 3(b)) throughout the whole temperature range. In $LaFeAsO_{0.9}F_{0.1}$, ^{19}F NMR showed small positive shift ~ 100 ppm above T_c [15]. This indicates that in $SmFeAsO_{0.84}F_{0.14}$, the dominant contribution to the Knight shift arises from the Sm-4f spins over that of Fe-3d. Moreover, in case of SmCoPO, where phosphorous lies in the Co-P plane, it was shown from ^{31}P NMR [16] that the major contribution to the Knight shift arises from Sm-4f electrons.

The T dependence of the shift, K is given by

$$K = K_0 + (H_{hf}/N\mu_B)\chi(T) \quad (1)$$

where K_0 is the T independent contribution arising from Pauli-paramagnetic and orbital parts of the conduction electrons, and transferred Van-Vleck susceptibility from rare-earth ion. H_{hf} represents the hyperfine coupling constant between the ^{19}F nuclear spin

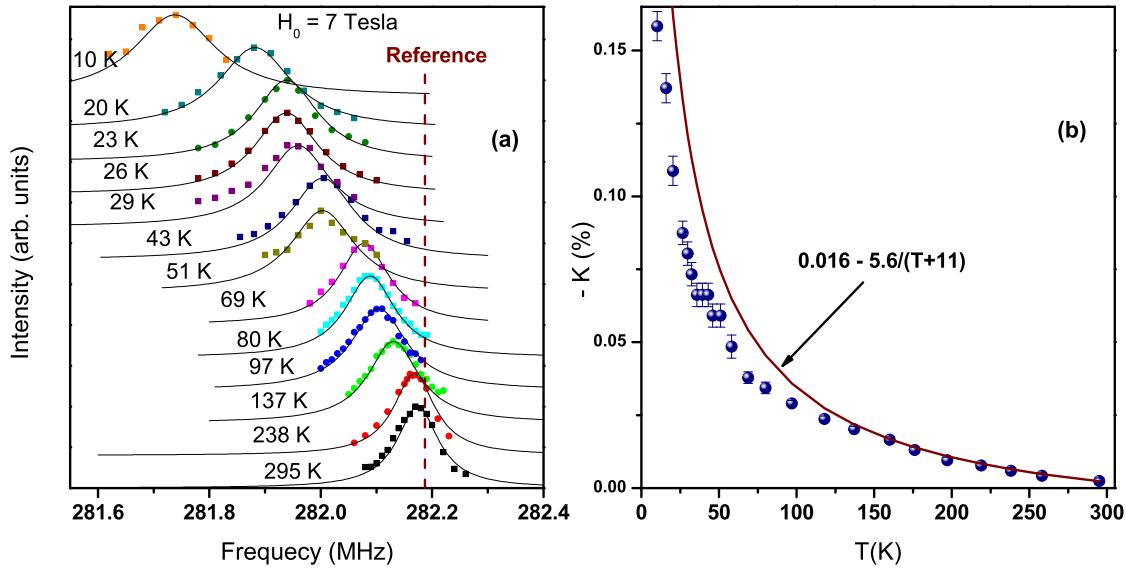


Figure 3. (Color online). (a): Frequency swept ^{19}F NMR spectra of $SmFeAsO_{0.86}F_{0.14}$ taken at 7.04 T. The continuous line represents the calculated spectrum. The vertical dashed line represents the resonance line position in a diamagnetic reference compound. (b): Temperature dependence of Knight shift (K) for ^{19}F NMR line in $SmFeAsO_{0.86}F_{0.14}$. Variation of $K(T)$ showing deviation from CW fit (—) below 160 K.

and the Fe-3d and Sm-4f electron spins, N and μ_B are the Avogadro number and Bohr magneton respectively. $K - T$ curve in the range 160 - 300 K, fitted (figure 3(b): solid line) using the CW formula for $\chi(T)$ gives $K_0 = 0.016\%$, $\theta = -11$ K, and the product of the C and H_{hf} . Using C from $\chi - T$ curve, we have $H_{hf} = -3.93 \pm 0.06$ kOe/ μ_B . The negative hyperfine coupling arises due to the dominant contribution of Sm-4f over that of Fe-3d, as the ^{19}F H_{hf} is positive in F-doped LaFeAsO [15]. In this connection it may be mentioned that the ^{19}F Knight shift data is not available in F-doped PrFeAsO.

3.3. Nuclear spin-spin ($1/T_2$) and spin-lattice $1/T_1$ relaxation rates

Figure 4 shows ^{19}F NMR transverse relaxation rate, $1/T_2$ (gives the intrinsic width of the resonance line) and the longitudinal relaxation rate, $1/T_1$ as a function of T . In the temperature range 10 - 300 K, except between a narrow range 160 - 170 K, T_1 was determined from the recovery of the longitudinal component of the nuclear magnetization $M(\tau)$ as a function of the delay time τ using $M(\tau) = M(\infty)(1 - \exp(-\tau/T_1))$. In the latter range we have used a stretched exponential ($M(\tau) = M(\infty)(1 - \exp(-\tau/T_1)^\beta)$) to fit the recovery curve with the exponent $\beta = 0.75$ (figure 5). The recovery of the transverse magnetization was found to be a single exponential at all temperatures. T_2 was obtained by fitting the decay curve with the

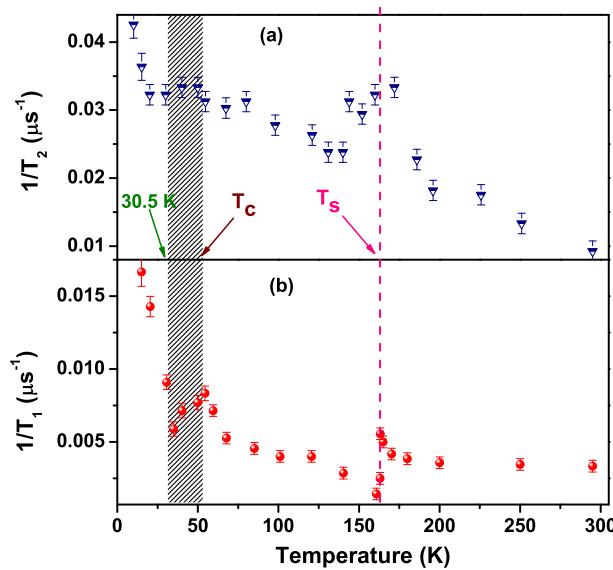


Figure 4. (Color online). Temperature dependence of ^{19}F NMR (a): spin-spin $1/T_2$ and (b): spin-lattice $1/T_1$ relaxation rate in $\text{SmFeAsO}_{0.86}\text{O}_{0.14}$.

equation $M(2\tau) = M_0 \exp(-2\tau/T_2)$, where M_0 is the initial magnetization.

Above 180 K, $1/T_1$ remains constant while $1/T_2$ shows continuous increment. Values of $1/T_1$ are almost same as reported by Prando *et al* [11] and three orders of magnitude larger than in $\text{LaFeAsO}_{1-x}\text{F}_x$ [15]. This implies dominant contribution of localized Sm $4f$ spin-fluctuation on the $1/T_1$ and $1/T_2$ processes [17, 18]. At around 163 K, both $1/T_1$ and $1/T_2$ data show a clear peak. This we associate to the SPT as reported by Martinelli *et al* [4]. We believe this observation is the first evidence of SPT in $\text{SmFeAsO}_{0.86}\text{O}_{0.14}$ using any microscopic tool. We shall discuss the nature of this SPT in the following section.

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Below SPT down to 50 K, both $1/T_1$ and $1/T_2$ show slow enhancement. However, in the range $30 \text{ K} < T < T_c$, $1/T_1$ shows a small but clear decrement which may be associated with the enhanced contribution from $4f$ local spin ordering superimposed on the decreasing trend of the diamagnetic contribution of the superconducting electrons. It is to be mentioned that the magnitude of ^{19}F NMR $1/T_1$ in the normal state of $\text{PrFeAsO}_{0.89}\text{F}_{0.11}$ is very close to that in $\text{LaFeAsO}_{0.89}\text{F}_{0.11}$, indicating a negligible

contribution of Pr-4f spin over Fe 3d [19]. Moreover, the effect of Pr-4f spin ordering on the T dependence of $1/T_1$ (as well as in Knight shift) was not observed. This makes it possible to observe a large drop in $1/T_1$ at T_c in $PrFeAsO_{0.89}F_{0.11}$.

Finally, the sharp increase of $1/T_1$ and $1/T_2$ in the range 10 - 25 K should be due to the development of Sm-4f spin short range correlation (which was also reflected in the behavior of the bulk susceptibility), as the system approaches T_N (3.8 K) [20].

3.4. The nature of the structural transition

We have shown the existence of anamoly (due to SPT?) near 160 K in the temperature variation of resistivity, susceptibility and the ^{19}F NMR Knight shift in $SmFeAsO_{0.86}O_{0.14}$. However, the most direct evidence of SPT has come from $1/T_1$ and $1/T_2$ measurement. Occurrence of SPT will be reflected from ^{19}F ($I = \frac{1}{2}$) NMR $1/T_1$ through a change in the electron-nuclear dipolar and hyperfine contribution to the fluctuating local magnetic field (H_{local}) produced at the ^{19}F site by the Fe-3d and Sm-4f spin. On the other hand $1/T_2$ is governed by the contributions from (i) dipolar interaction of the ^{19}F nuclear spins with the like and unlike nuclear spins, which is temperature and field independent ($1/T_2|_{static}$) and (ii) the dipolar and hyperfine interactions of the ^{19}F nuclei with the longitudinal component of the fluctuating magnetic field produced by the neighboring Fe^{2+} 3d-spins and Sm 4f-spins ($1/T_2|_{dynamic}$). For localized spin, the latter contribution shows a Curie-Weiss behavior in T dependence [18]. Close to the magnetic ordering temperature, it shows strong T dependence, as the fluctuation frequency of H_{local} reduces and becomes comparable to the nuclear resonance frequency due to the development of short range magnetic correlation. Near SPT (T_S), as the atomic positions are altered, it should change both ($1/T_2|_{static}$) and $1/T_2|_{dynamic}$. The magnitude of the drop in both $1/T_1$ and $1/T_2$ (clearly revealed from the change in the nature of the decay curves (figure 5)) below T_S is small indicating a weak tetragonal to orthorhombic distortion as also pointed out by Martinelli et al from structural study. Nevertheless, the effect of this small orthorhombic distortion is reflected in both $1/T_1$ and $1/T_2$ data. As pointed out earlier, the T independent behavior of $1/T_1$ in the range 180 - 300 K indicates faster fluctuations of H_{local} than the NMR resonance frequency (ν_R). The enhancement of $1/T_1$ in the range 163-175 K is a clear signature of the slowing down of the fluctuation of H_{local} , so that its frequency is comparable to ν_R , contributing to the relaxation process. Such slowing down of the fluctuation of H_{local} can arise due to several reasons. (1) It may be due to the development of a short range magnetic correlations of Fe-3d spins (as the Sm-4f spins order magnetically at 3.8 K (T_N), the short range correlation among the 4f spins would occur close to T_N , so its temperature dependent contribution to $1/T_1$ would arise far below 175 K). This is supported by the fact that χ and K vs T curve deviate from CW law below 160 K (figure 2(b) and 3(b)) with $\theta = -11$ K implying AFM correlations. Thus a role of AFM correlated spin fluctuations (may be secondary) in driving the SPT can not be ruled out. (2) The effect of slowing down of the orbital fluctuations near T_S may also have a contribution (through spin-orbit coupling), if SPT

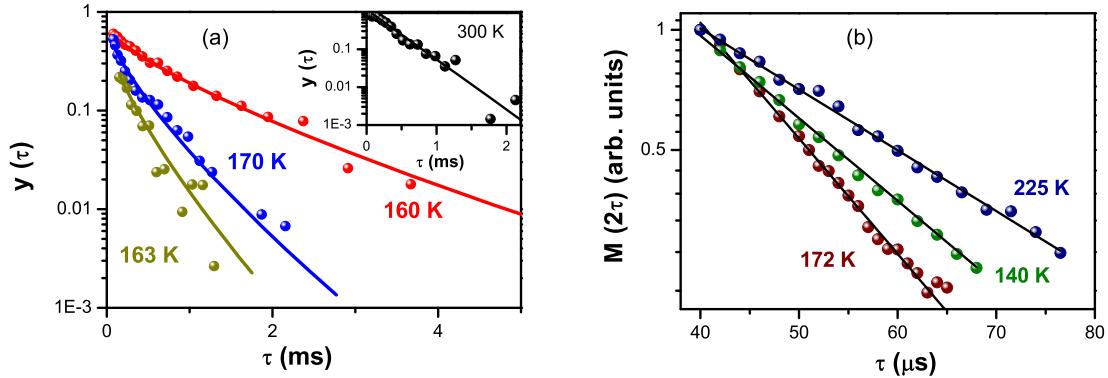


Figure 5. (Color online). Recovery curves for (a): longitudinal (T_1), ($y(\tau) = \frac{M(\infty) - M(\tau)}{M(\infty)}$) and (b): transverse (T_2) relaxation times of $\text{SmFeAsO}_{0.86}\text{F}_{0.14}$.

is driven by $3d$ orbital ordering as proposed theoretically [21] and from synchrotron powder diffraction data [4]. (3) Softening of a lattice vibrational mode near T_S can also enhance $1/T_1$ provided it can reduce the fluctuation frequency of H_{local} through a coupling with the spin /orbital motions of the electrons, as the contribution of phonon modes to $1/T_1$ is much less compared to that of the fluctuations of H_{local} .

We have shown that within a short interval 170 - 160 K, the recovery curves of T_1 (figure 5(a)) follow a stretched exponential $M(\tau) = M(\infty)(1 - \exp(-\tau/T_1)^\beta)$ with the exponent $\beta \sim 0.75 \pm 0.03$ indicating a distribution of $1/T_1$. Spatial distribution of the SPT temperature can induce a stretched exponential in the recovery curve but the clear sharp peak in $1/T_1$ (figure 4(b)) eliminates this possibility. However, the SPT related disorder may induce a distribution of T_1 over a certain temperature range above T_S resulting in the necessity of a stretched exponential. At the final stage of preparation of this manuscript, we have noticed the work of Martinelli *et al* [22], where they have shown that on cooling, micro strain along the tetragonal $hh0$ direction appears and increases as the temperature is decreased. Just above the structural transition, micro strain reaches its maximum value and then is abruptly suppressed by symmetry breaking. Micro strain reflects a distribution of the lattice parameters in the tetragonal phase and hence explain the occurrence of stretched exponential in the recovery curves of T_1 . From NMR results it is therefore confirmed microscopically that in optimally F doped SmFeAsO system which exhibit superconductivity, the tetragonal to orthorhombic transition near 163 K is still present, as in the non superconducting parent compound. This finding creates interest to explore whether this occurs only in Sm based Fe-As systems or also in other rare-earth based 1111 superconductors.

4. Conclusion

In conclusion, we report the resistivity, magnetization and ^{19}F NMR results in superconducting $\text{SmFeAsO}_{1-x}\text{F}_{14}$. Both the resistivity and the magnetization data clearly show the signature of the superconducting transition at 48 K. The behavior

of $1/T_1$ and $1/T_2$ confirms, for the first time from a microscopic tool like NMR, the existence of a structural phase transition at 163 K, as in the non superconducting parent compound, in a sample which undergoes superconducting transition at 48 K. This finding creates interest to explore whether this is unique for only Sm based systems or it is also present in other rare-earth based 1111 superconductors.

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